

# Constructing your very own DIY dimensional world-line portal; measuring and creating frequency profiles of location (part 3)

This post continues in the discussion of building yourself a DIY dimensional portal (or some type of vehicle) for world-line crossovers and slides. This is part three. Part one was an introduction to the concepts that various people can build a DIY dimensional portal. Part two discussed the very important aspects of mass / gravity separation of the entity (person) entering the portal, and the portal itself.

And here, in this part we will discuss measuring the frequencies of the gravity elements involved when a person enters the portal. This measurement of frequencies is the assignment of coordinates of where you are right now at the moment of teleportation.

Measure frequencies = Assign egress coordinates.

## High Frequency Gravity Waves

The fundamental idea is that we would detect the super weak HFGW that is associated with both the mass of the person entering the portal, and

that of the portal itself. This would create a frequency profile. This profile in turn, can be considered a set of coordinates for the dimensional portal to work with.

Gravitational waves (GW) are a prediction of Einstein's general theory of relativity, but (due to their weakness) took a long, long time to discover.

Measurement of their indirect effects on the orbits of certain binary neutron stars was a major experimental triumph, and merited the award of a Nobel Prize in Physics. Further; these measurements agree with theory to better than 1%. Therefore, there really isn't any question of their existence. The issue is really how to detect them for small gravitational masses, up close, quickly and accurately.

The term HFGW has come to mean gravitational waves at much higher frequencies of several GHz, say 10GHz to be specific. A general rule of thumb is that the *highest* gravitational wave frequencies produced will be at around the reciprocal of the freefall timescale for a system  $f_{\max} \sim \sqrt{G\rho}$ , where  $\rho$  is the average density of the system.

Dr. Robert Baker, Jr. has a design for an open cavity High-Frequency Gravitational Wave Detector in the GHz band. His design consists of a high-quality-factor open microwave cavity and a Gaussian beam (GB) passing through a static magnetic field in free space.

Baker is regarded as the preeminent researcher in the field of High-Frequency Gravitational Wave research, and proposes this new detector model as a means of facilitating significant new potential applications for the wireless telecommunications sector.

Essentially this effect is an inverse Gertsenshtein effect in which HFGWs are converted into electromagnetic (EM) waves when passing through a static magnetic field.

Our dimensional portal would detect the isolated HFGW's from both the portal and the person entering the portal. It would convert the values into electromagnetic waves when the person enters the dimensional portal. Of course, for this to work, the entire portal would need to be a static magnetic field.

## The Physics of HFGW's

Newton's formulation of the theory of gravity,

$$\vec{F}_{12} = -\frac{M_G(1) M_G(2)}{r_{12}^2} \left( \frac{\vec{r}_{12}}{r_{12}} \right) \quad (3-1)$$

$$\vec{F} = M_I \vec{a} \quad (3-2)$$

$$M_I = M_G \quad (3-3)$$

for two spherical gravitating masses  $M_G(1)$  and  $M_G(2)$  is equivalent to the "non-relativistic" gravitational field description

$$\vec{\nabla} \cdot (\vec{\nabla} \hat{h}) = \frac{4\pi G}{c^2} \rho_G \quad (3-4)$$

$$c^2 \rho_G \vec{\Delta} \hat{h} = \rho_I \vec{a} \quad (3-5)$$

$$\rho_I = \rho_G \quad (3-6)$$

in which a non-dimensional “potential”  $\hat{h}$  has been chosen to agree with the mathematical language used for it in General Relativity. Here  $M_I$  and  $M_G$  are the inertial and gravitational masses respectively, and  $\rho_I$  and  $\rho_G$  are the distributions of these masses. Equations (3-4) and (3-5) are an instantaneous action-at-a-distance description which is inconsistent with the constraints of Special Relativity.

In General Relativity (which generalizes Newton’s theory) Equations (3-4) – (3-6) become

$$\left( \nabla^2 - \frac{\partial^2}{c^2 \partial t^2} \right) \hat{h}_{\mu\nu} = \frac{4\pi G}{c^2} \hat{T}_{\mu\nu} \quad (3-7)$$

with

$$\hat{T}_{\mu\nu} = 2 \left[ T_{\mu\nu} - \frac{g_{\mu\nu}}{2} T \right] / c^2 \quad (3-8)$$

$T_{\mu\nu}$  is the complete relativistic stress-energy tensor of everything including the gravitational field itself, and  $T$  is its trace. ( $g_{\mu\nu}$  is the Minkowski metric tensor of Special Relativity plus  $\hat{h}_{\mu\nu}$ .) Confirmed predictions include the equivalence principle  $\rho_I = \rho_G$  (to better than  $10^{-10}$ ), the calculated value for the bending of light passing near the sun and gravitational lensing of light in other parts of the Universe, many solar system observations, and remarkably accurate observations of neutron star binaries.

The full content and implications of General Relativity are not needed for any of the HFGW predictions to be considered below. For example the quantum energy density in a vacuum is negligibly small compared to the other important matter and field contributions to  $T_{\mu\nu}$  in our local environment. All of the HFGW amplitudes of interest here are so small that their contributions to energy density can be neglected in  $T^{\mu\nu}$ .

In a vacuum with only  $h^{\mu\nu}$  present the RHS of Equation (3-7) vanishes, leaving the familiar free field wave equation

$$\left(\Delta^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2}\right) \hat{h}_{\mu\nu} = 0 \quad (3-9)$$

The robustness of the basic theory for the HFGWs discussed below is even more robust than that of General Relativity.

Hypotheses about changes in gravity and  $T_{\mu\nu}$  from string theory might change it at length scales  $\ll 1$  cm and some have proposed changes at huge (astronomical/cosmological) scales but neither would change Equations (3-7) on the scales of interest here.

Because we are concerned with such small HFGW intensities it is often constructive to describe these flows as a flow of gravitational quanta (gravitons).

Gravitons are a necessary consequence of Quantum Mechanics applied to Equation (3-9) and bear the same necessary relationship to Equations (3-9) and (3-7) as photons do to electromagnetic fields.

In particular

$$E \text{ (graviton)} = \hbar \omega = \hbar c k, \quad (3-10)$$

with  $\omega = 2\pi \times$  frequency and  $k = 2\pi/\lambda$ .

Figure 1 shows the electromagnetic-gravity field interactions in Equation (3-7) as (static gravity or graviton) – (photon or static electromagnetic field) interactions.

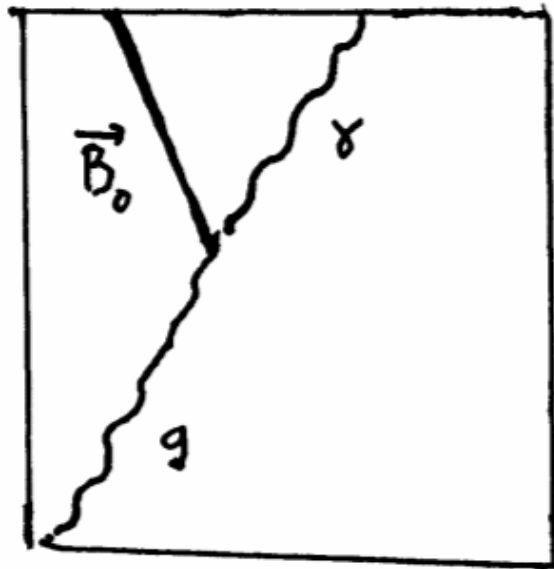
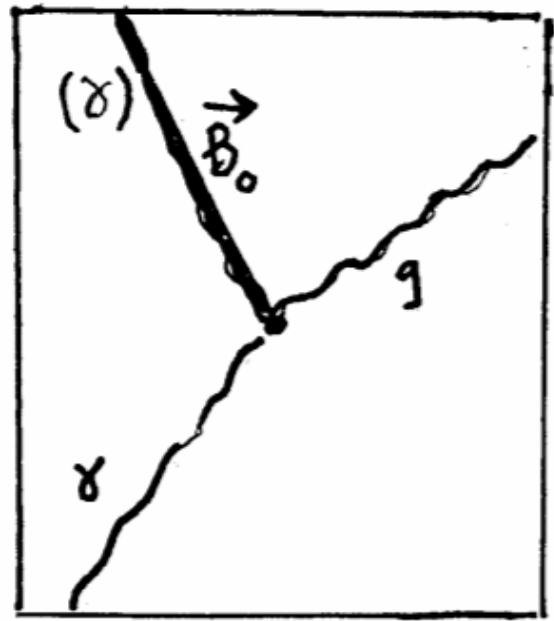
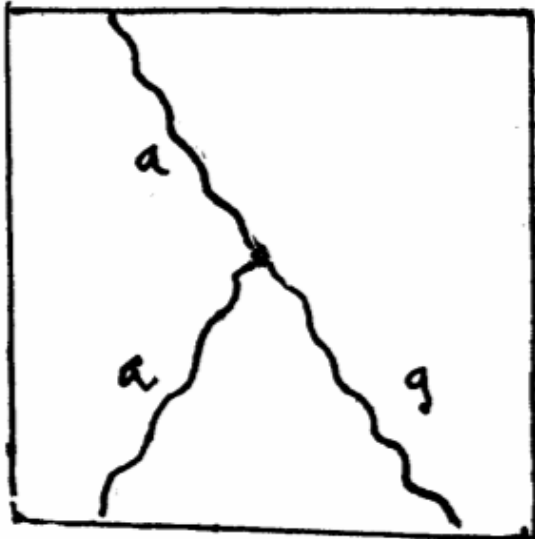
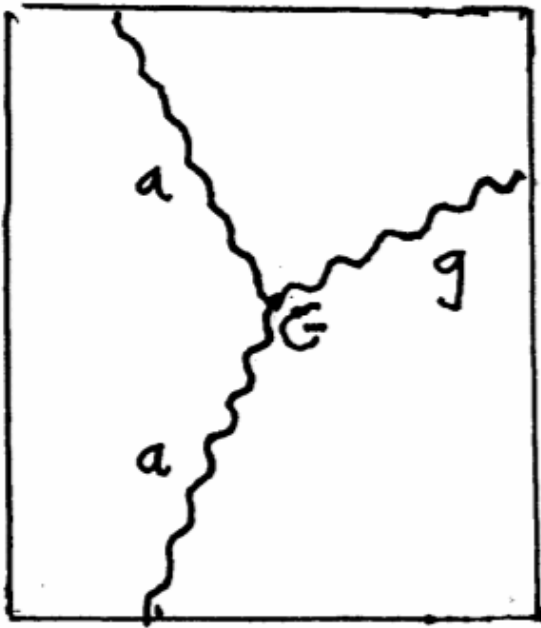


Figure 1: Feynman diagrams of quantum (graviton/photon) reactions in quantized gravitational field versions of General (and Special) Relativity.  $\gamma \equiv$  HF electromagnetic field or static field ( $B \neq 0$ );  $g \equiv$  graviton;  $A \equiv$  any particle.

# Measuring HFGW from gravity masses

The LIGO detectors, which measured the waves, do not use bar detectors; they use interferometers. Bar detectors have been used for decades, but they have not been sensitive enough to make actual detections. They are necessarily very short, which reduces the effect of a gravitational wave. As you indicate they also have fairly narrow resonant frequencies at which they are most sensitive. Interferometers, on the other hand, can be made 4 kilometers long (like the LIGO detectors), which magnifies the effect of the waves. They are also sensitive over a fairly broad range -- roughly 40Hz to 2000Hz.

As anna v rightly points out, there actually are plenty of references to frequency if you look at the science papers. I work in gravitational-wave astronomy, and decomposing things into frequencies is our bread and butter. There's less coverage of this in the popular press, presumably because the public tunes out talk of frequencies, and pop-sci journalists know where their bread is buttered. But Fourier transforms are really how the analysis gets done.

-Physics Stack Exchange

Dr. Robert Baker, Jr. has a design for an open cavity High-Frequency Gravitational Wave Detector in the GHz band, which consists of a high-quality--factor open microwave cavity and a Gaussian beam (GB) passing through a static magnetic field in free space.

Essentially this effect is an inverse Gertsenshtein effect in which HFGWs are converted into electromagnetic (EM) waves when passing through a static magnetic field.



# Converting measured HFGW into electromagnetic waves for frequency generation.

A basic mechanism for generating a EM wave from a measured HFGW is the direct conversion of the same frequency by a strong static magnetic field ( $\rightarrow B_0$ ).

This Gertsenshtein process is idealized in Figure 3. The GW power out,  $P_{GW}$  (out), is proportional to the electromagnetic wave incoming power  $P_{EMW}$  (in):

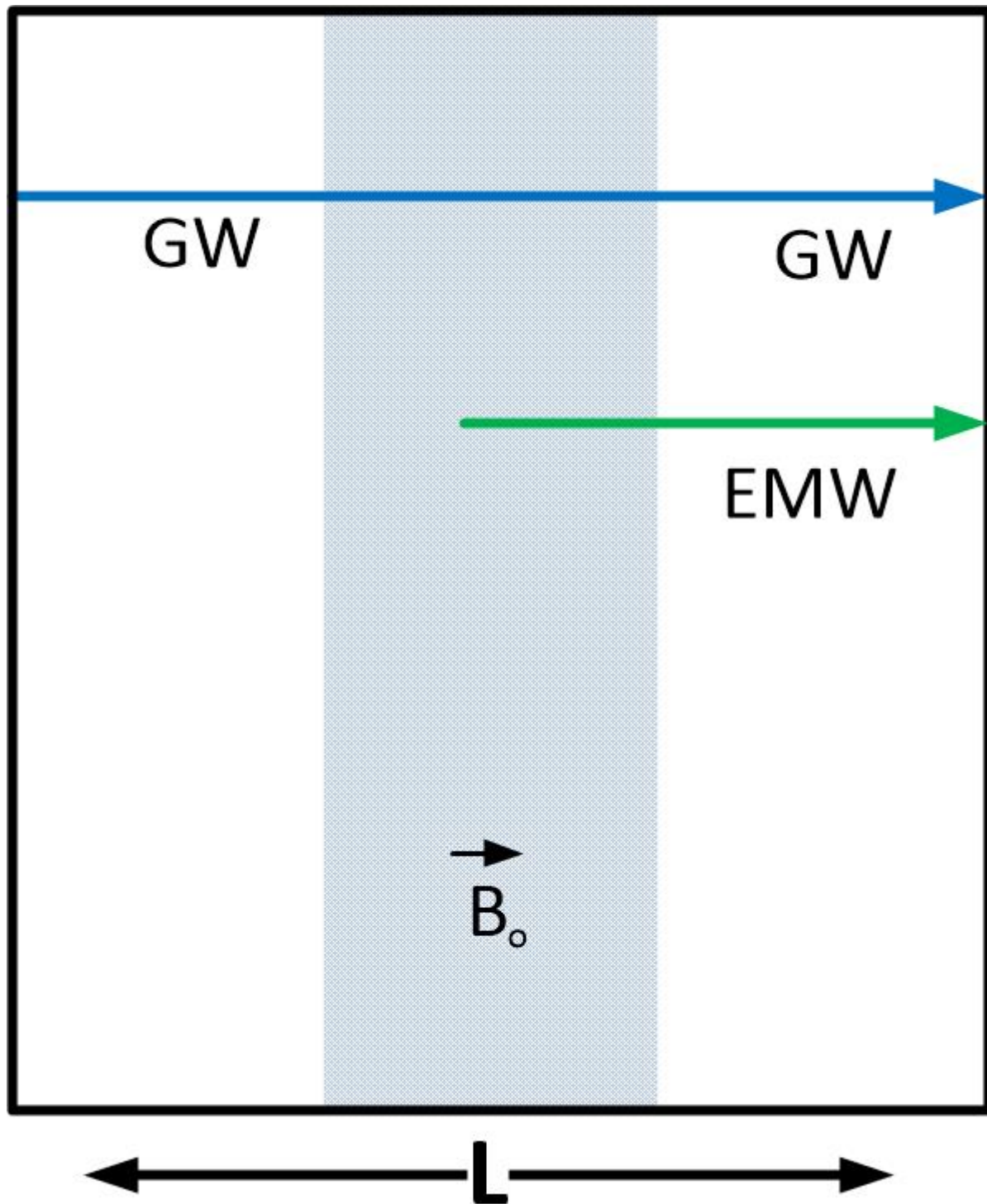


Figure 3: Gertsenshtein EMW generation by HFGWs passing through a constant magnetic field  $B_0$ ,

$$P_{GW} \text{ (out)} = \mathcal{F} P_{\text{EMW}} \text{ (in)} \quad (3-13)$$

$$\mathcal{F} = \frac{4 \pi G B_0^2 L^2}{c^4} \sim 10^{-35} \quad (3-14)$$

for  $B_0 = 10^5$  Gauss and  $L^2 = 10^3 \text{ cm}^2$ . Equivalently

$$P_{GW} \text{ (out)} = \frac{4 \pi G B_0^2 U}{c^3} L, \quad (3-15)$$

where  $U$  is the total EMW energy in the volume ( $V$ ) in which the EMW passes through  $B_0$ .

$$u = \frac{U}{V}$$

is the energy density in that region.

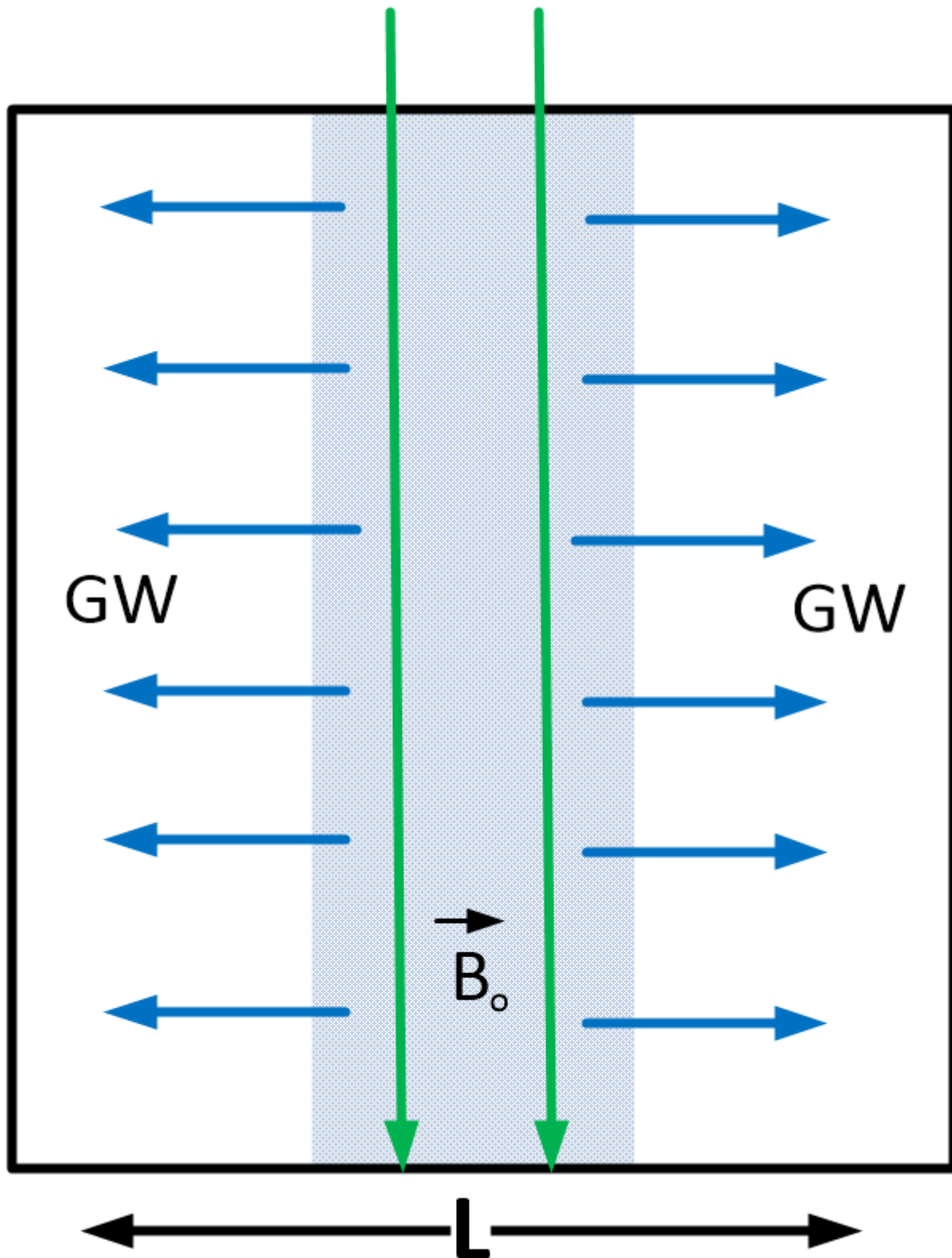


Figure 4: HFGW generation by standing wave electromagnetic modes in a cavity.

For the geometry of Figure (3) in which the passage of the EMW through  $B_0$  is not otherwise interrupted

$$U = P_{\text{EMW}} (\text{in}) L/c. \quad (3-16)$$

For  $P(\text{in}) \sim 10 \text{ kW}$ , and  $L = 30 \text{ cm}$ ,  $U = 10^{-5} \text{ joules}$ . If the EMF is contained as a normal mode within  $V$ ,  $U$  can be very much larger. However, there are various limits to  $U$  which are independent of the available EMW power. For a cavity with EM dissipation time  $\tau$

$$\omega \tau \equiv Q \quad (3-17)$$

the heat loss rate

$$\dot{H} = \frac{U}{\tau} = \frac{U \omega}{Q} \quad (3-18)$$

For a (generous) cooling rate from an exterior coolant flow around a copper cavity  $H' \sim 10^6 \text{ watts}$ ,  $Q \sim 2 \times 10^3$ ,  $U_{\text{max}} \sim 2 \times 10^{-1} \text{ joules}$  and

$$\text{Max } P_{\text{GW}} (\text{out}) \sim 2 \times 10^{-27} \text{ watts} \sim 2 \times 10^{-2} \text{ graviton/sec.} \quad (3-19)$$

(We note that it would take a continual EM power input of one MWatt to maintain this tiny GW output.)

If we replace the copper-walled cavity by one with superconducting walls  $\tau$  may increase from the  $\sim 10^{-7} \text{ sec}$  of Cu by a factor  $\sim 10^7$ . However,  $U_{\text{max}}$  could not increase by nearly such a factor, even if we ignore any problems of maintaining superconductivity near the huge  $\rightarrow B_0$ , and keeping the very

low temperature needed. The  $u$  inside the superconducting cavity would be limited by unacceptable electron emission from a mode's strong electric field perpendicular to a wall:

$$E_{\perp} < \frac{50 \times 10^6 \text{ volts}}{\text{meter}} \quad (3-20)$$

This implies

$$u \sim E_{\perp}^2 / 8\pi < 10^7 \text{ erg cm}^{-3}. \quad (3-21)$$

Then with an assumed  $V \sim 3 \times 10^3 \text{ cm}^3$

$$\text{Max } P_{\text{GW}} (\text{out}) \sim 10^{-23} \text{ watts} \sim 10^2 \text{ graviton/sec} \quad (3-22)$$

Even if this crucial limit is ignored there would be a limit to  $u$  from the maximum mechanical strength of the container confining the electromagnetic modes:

$$u_{\text{max}} \sim 10^{10} \text{ dyne/cm}^2. \quad (3-23)$$

The limit of Equation (3-23) and  $V \sim 3 \times 10^3 \text{ cm}^3$  gives  $U_{\text{Max}} \sim 3 \times 10^6 \text{ J}$  and

$$P_{\text{GW}} (\text{out}) \sim 10^{-20} \text{ watts} \sim 10^5 \text{ graviton/sec} \quad (3-24)$$

Finally we could ask the ultimate limit when, instead of  $\rightarrow B_0 \sim 10^5 \text{ Gauss}$

and EM waves  $V$  is filled with moving masses, EM energy, etc. all contained within  $V \sim 3 \times 10^3 \text{ cm}^3$  to the limit where the container explodes. Then

$$P_{GW} \text{ (out)} \sim 10^{-18} \text{ watts} \sim 10^7 \text{ graviton/sec.} \quad (3-25)$$

The graviton flow at a target a distance  $d$  away is

$$\frac{P_{GW} \text{ (out)} \times b}{4 \pi d^2} \quad (3-26)$$

where  $d$  is the distance to the target and  $b$  a directional beaming factor which we take  $\sim 10^2$ . Then for  $d > 1 \text{ km}$  the maximum flux at a target

$$f < 10^{-9} \times (10^7) \frac{\text{graviton}}{\text{cm}^2 \text{ sec}} \quad (3-27)$$

for the unrealistically large limit of Equation (3-25). Increasing  $V$  to  $10^7 \text{ cm}^3$  would still limit

$$f < 30 \text{ gravitons/cm}^2/\text{sec.} \quad (3-28)$$

Almost none will be stopped or converted within the target. (But even

if they were their total impulse would cause no damage to any part of it.)

## HFGW Detectors [1]

Proposed HFGW detectors have generally been based upon versions of the inverse Gertsenshtein process. The most elementary one is that in Figure 5. As in Equations (3-13) and (3-14)

$$P_{EM} \text{ (out)} = \mathcal{F} P_{GW} \text{ (in)} \quad (3-29)$$

and

$$\mathcal{F} = 10^{-35} \text{ for } L^2 = 10^3 \text{ cm}^2 \text{ and } B_0 = 10^5 \text{ Gauss.}$$

For the maximum HFGW generator production of  $10^2$  graviton/sec of Equation (3-22), and  $b \sim 10^2$  and  $d \sim 10$  m in Equation (3-26), and a detector area transverse to the beam  $(A^{\wedge}) = 10^4 \text{ cm}^2$

$$P_{EM} \text{ (out)} \sim 10^{-34} \text{ EM photons/sec} \equiv \dot{N}_{\gamma}$$

with an average interval between photons

$$\hat{t}_{\gamma} = \frac{1}{\dot{N}_{\gamma}} = 10^{34} \text{ sec.} \quad (3-30)$$



Such a small photon flow would, of course, never be observed, no matter what plausible changes are made in HFGW generator,  $d$ ,  $b$ , or  $A^{\wedge}$ . However proposals have been made to decrease this interval by very great factors.

One such proposal introduces an additional EMW $\theta$  with the same frequency as the GW and the very weak EMW it generates in passing through the strong  $\rightarrow B\theta$  region. This is well understood "homodyning" of the weak signal. It does not increase a signal to noise ratio when the noise is the minimal photon noise from quantization. If we consider the simple geometry of

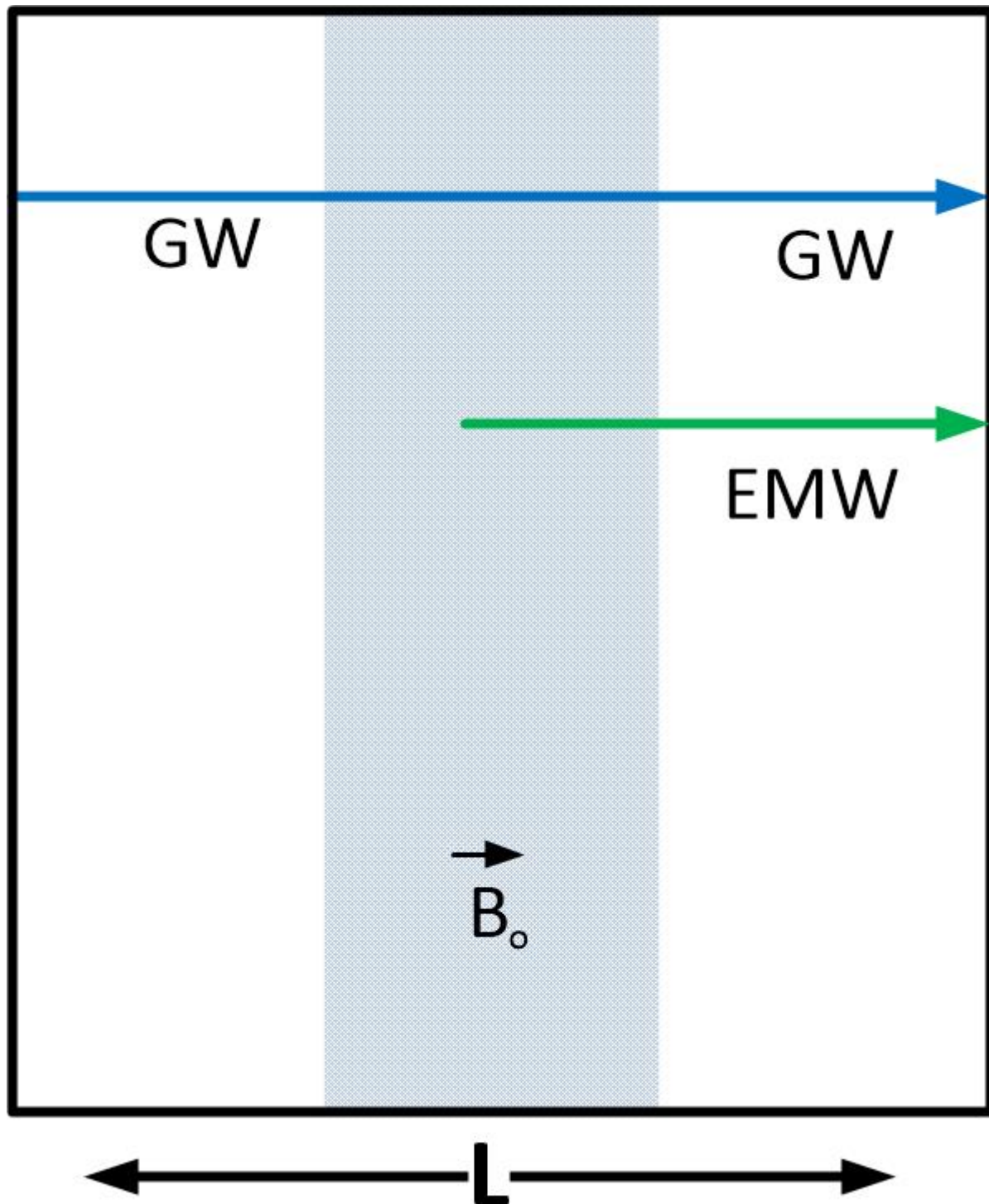


Figure 5: Inverse Gertsenshtein conversion of HFGWs to EMWs of the same frequencies.

Figure 6 with the electromagnetic waves electric field normal to the plane of wave propagation and  $\vec{B}_0$ , there are two possibilities for interference between EGW, the electric field of the EMW generated by the GW and  $E_0$ . In one the original propagation directions are coincident. Then the total field ( $\vec{E}_T$ )

$$\vec{E}_T = \vec{E}_0 + \vec{E}_{GW},$$

with  $\vec{E}_T$  the homodyning field and  $\vec{E}_{GW}$  that from GW conversion along the common trajectory. If  $\vec{E}_{GW}$  reaches the photon detector so must  $\vec{E}_T \neq 0$ . That detector's photon counting rate

$$\dot{N}_d \propto |\vec{E}_T|^2 = (|\vec{E}_0|^2 + 2\vec{E}_0 \cdot \vec{E}_{GW} + |\vec{E}_{GW}|^2) \quad (3-31)$$

After a long time  $t$  the collected number of photons

$$N_d = \dot{N}_0 t + 2(\dot{N}_0 \dot{N}_{GW})^{1/2} t \cos \delta + \dot{N}_{GW} t \quad (3-32)$$

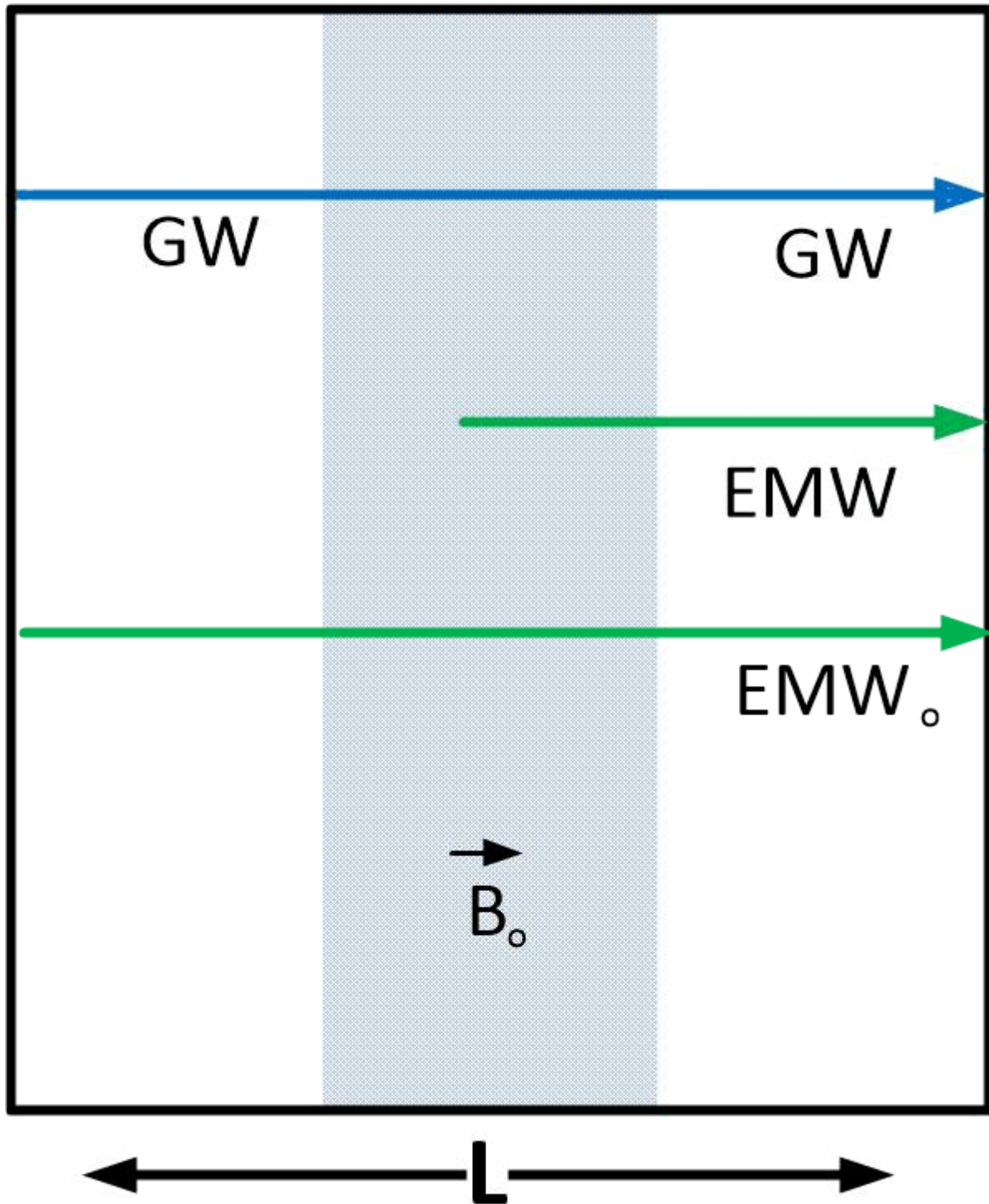


Figure 6: Homodyning of weak EMW with much stronger EMW<sub>0</sub>.

with  $N'_{\theta}$  the counting rate when  $N'_{GW} = 0$  and  $N'_{GW}$  the very much smaller rate when  $N'_{\theta} = 0$ . A non-zero  $\cos \delta$  can arise from phase match between  $\rightarrow E_{\theta}$  and  $\rightarrow E_{GW}$ .

The large  $N_0 = N'_{\theta t}$  is the expectation value of a Poisson distribution

of width  $N^{1/2} \theta$  which is intrinsic to the quantum (photon) distribution in the classical wave description.

The main  $N \dot{N}_{GW}$  contribution to the detector counts  $(2 (N \dot{N}_{GW})^{1/2} \cos \delta t)$  must be significantly larger than this fluctuation  $(N \dot{N}_0 t)^{1/2}$  for the signal/minimal photon noise ratio to exceed unity:

$$(\dot{N}_0 \dot{N}_{GW})^{1/2} t > (\dot{N}_0 t)^{1/2}. \quad (3-33)$$

Then

$$t > 1/\dot{N}_{GW} = \hat{t}_\gamma, \quad (3-34)$$

i.e., it will still take the  $\hat{t}_\gamma$  of Equation (3-30) to identify with any confidence a single EMW photon from incoming GW graviton conversion.

If the  $\rightarrow E \theta$  photons differ enough in direction from the  $E \leftarrow GW$  ones so that they do not reach the detector the photon fluctuations  $|\rightarrow E \theta|^2$  term of Equation (3-31) could be absent, but so would  $2 \rightarrow E \theta \cdot \rightarrow E GW$  so that again  $t \sim 1/N \dot{N}_\gamma$ . The history of this interference term before the detector is reached is not relevant:  $t \sim 1/N \dot{N}_{GW}$  whether or not  $\rightarrow E \theta$  reaches the photon detector with  $\rightarrow E GW$  or what its magnitude there is as long as it gives the minimal fluctuation in photon number as the major noise source at the EMW detector.

If instead of  $\rightarrow E \theta$  with the same frequency at the EMW from HFGW conversion (homodyning), the  $\rightarrow E \theta$  wave has a different frequency ( $\omega$ ) and the detector admits  $\omega \pm \omega$  (heterodyning) the quantum limit still gives the same needed  $t$  (to within a factor 2) for a signal to noise ratio exceeding one; see Marcuse [13] (Eqs. 6.5-14, 6.5-17) with the minimum band-

width  $B \sim t^{-1}$  achieved over a time  $t$ ,

## HFGW Detectors [2]

A second kind of proposal for greatly increasing the photon counting rate from graviton  $\rightarrow$  photon conversion is to contain the conversion volume within reflecting walls for EMWs.

This is similar to the same sort of proposal to increase the efficiency of Gertsenshtein conversion of photons to gravitons in Figure 3. It differs, however, in that the containing cavity does not reflect the gravitons which are the source for conversion, but only the photons which are the product of it.

If we start with an empty cavity with mode decay time  $\tau$  and a resonance frequency  $\omega_0 = \omega$  (or at least  $|\omega - \omega_0| < \omega_0/Q$ ) the cavity will initially fill with EM mode energy ( $U$ ) at a rate

$$\dot{U} \sim \frac{\hbar \omega_0}{\hat{t}_\gamma} \left( \frac{ct}{L} \right) \quad (3-35)$$

which will continually increase until a steady state is reached at  $t \sim \tau \equiv Q/\omega$ . ( $U$  is not limited in the cavity detector by the considerations of Sec 3.

because it is always so tiny in comparison to those in a GW generator).

At  $t = \tau$  a maximum photon counting rate

$$\dot{N}_\gamma \sim \frac{1}{\hat{t}_\gamma} \left( \frac{c \tau}{L} \right) = \frac{Q}{\hat{t}_\gamma} \quad (3-36)$$

if cavity photons are counted instead of being dissipated in the cavity walls.

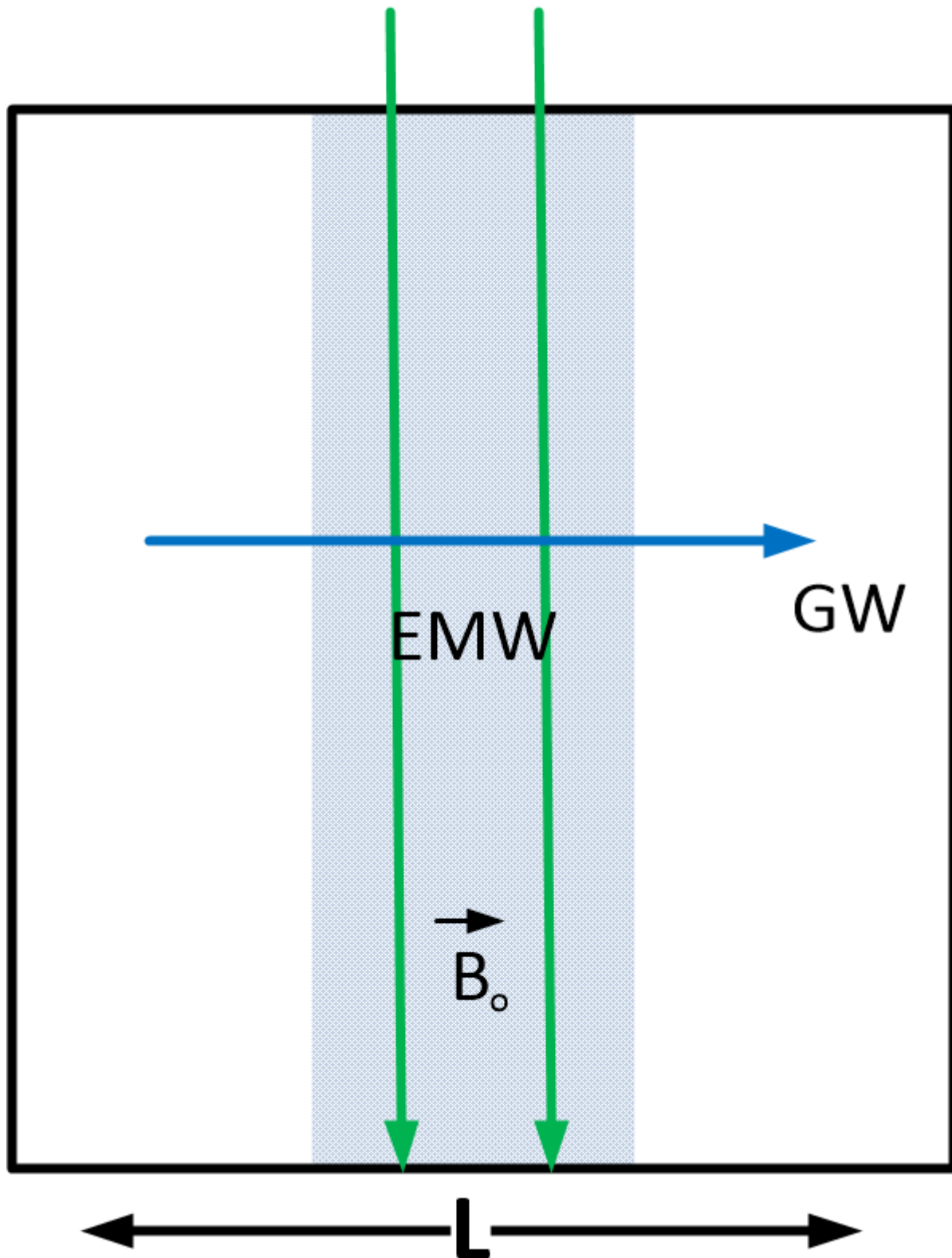


Figure 7: GW conversion on  $B_0$  pumping a resonant cavity with the same frequency.

If, unphysically, finite cavity mode decay time did not limit  $N \gamma$  we might still note how long ( $t_1$ ) it would take for the expected number of GW induced photons inside the cavity to reach one, i.e.



$$1 = \frac{U}{\hbar \omega_0} = \left( \frac{c t_1}{L} \right) \frac{t_1}{\hat{t}_\gamma}. \quad (3-37)$$

Then

$$t_1 \sim \left[ \left( \frac{L}{c} \right) \hat{t}_\gamma \right]^{1/2} \sim 10^{12} \text{ sec.} \quad (3-38)$$

However, finite  $\tau = (Q/\omega)$  does limit the cavity  $U$ . The maximum expected value for GW induced photon number in the cavity never approaches unity. Instead

$$\frac{U_{\max}}{\hbar \omega_0} \sim \left( \frac{\tau}{\hat{t}_\gamma} \right)^2 \ll 1 \quad (3-39)$$

A copper-walled cavity with  $Q \sim 2 \times 10^3$  would decrease the time interval between GW induced photons in the cavity, but only to

$$\hat{t}_\gamma / Q \sim 10^{31} \text{ sec.} \quad (3-40)$$

The largest plausible  $\tau$  would be for a cavity with superconducting walls. Then  $\tau$  might reach, say, 10 seconds ( $Q \sim 10^{11}$ ). Then

$$\hat{t}_\gamma / Q \sim 10^{23} \text{ sec,} \quad (3-41)$$

still essentially an infinite time between photon counts.

If the cavity GW induced photon energy were homodyned (or heterodyned) by introducing additional resonant mode electromagnetic field energy the photon number fluctuations in that energy would again not allow interference to increase the time interval for signal/photon noise  $> 1$  to be less than the  $\hbar\omega/Q$  of Equations (3-40)- (41).

## What this means

There is a way (of a couple of ways) to measure the gravity waves associated with the gravity of a person entering a portal, and that of the portal itself. These waves at a precise moment in time can be used as a coordinate.

It is not practical to use this technology for any other purposes.

The photon counting rates for confident detection of graviton-induced photons from proposed HFGW generators and detectors is so small that development of HFGW communication links is not a reasonable prospect.

- Not useful for communication.

The graviton interception-transformation rate at a large cooperative target (specially designed to detect gravitons)  $\approx 10^{-20}$  [ cf Equations (3-29)

and (3-36)]. When combined with the comparably small fraction for photo  $\rightarrow$  graviton efficiency in HFGW generators this implies that to deposit even an

ergs worth of HFGW gravitons in a target requires  $\approx 10^{40}$  ergs of electric power input to a HFGW generator. This is more than total energy from electric power generation on the earth ( $< 10^{12}$  watts) for longer than the age of the Universe.

Use of HFGW beams for destroying, deflecting, or compromising distant targets (or close ones) has no promise.

- Not useful for weapons.

Thus it seems silly that the United States government would consider putting this technology in a “black project” to keep it out of the public eye.

## Conclusion

This part discussed creation of a mechanism to measure the gravity waves associated with the gravity of both the dimensional portal and a person entering it.

With this mechanism you can identify the exact world-line you are in at an exact frozen moment of time, and assign a coordinate to it.

You can do so in isolation of the person, and thus create a mechanism that would take this “person” at one coordinate and slide him to another coordi-

nate instantaneously.

Since the coordinate is very detailed, it includes not only the physical geography of a place, but a moment in “time”, and if you change the coordinates slightly, you can use this mechanism to move a person back and forth in ...

- Geography. You can move about from place A to place B.
- Time. You can move from one point in time to another.

But since, you have the entire spectrum of coordinates at your “finger tips” you can alter the parameters of the coordinates to enter completely different world-lines. You can go into the so-called parallel universes...

- World-line. You can go from one world-line to another.

In the next post, we will discuss how to use these frequencies to move a person from one set of coordinates to another set. Hang on...

## Do you want some more?

I have more posts on this subject here...

---

I have more posts of a similar nature in my MAJestic index [here...](#)